Polarized Tops from New Physics

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Tops and new physics at the LHC

- Top quark physics a highly anticipated aspect of physics at the LHC
 - o top factory: SM $t\bar{t}$ production $\sim 830~\mathrm{pb}$
- New physics responsible for EWSB must have large couplings to the top
 - Anticipate relatively large cross-section for associated production of top and new physics
- New physics will couple differently to t_L , t_R
 - $\circ \Rightarrow$ net polarization of tops from new physics
 - o polarization will depend on kinematics and chiral couplings
- Polarization of tops from new physics probes chiral structure of new physics

Polarization without reconstruction

- Approach: study top polarization using kinematic distributions which do not require full event reconstruction
- With missing energy, event reconstruction requires assumptions about event topology
 - Not necessarily desirable, especially in initial phases: new physics with stable invisible particles, complicated cascade decays
- Hadronic tops: top can be completely reconstructed without information from the rest of the event
- Collimated tops: can construct variables which become independent of the top rest frame in the collinear limit
 - works for *both* hadronic and leptonic tops

Hadronic tops

• Angular distribution of daughter particles in top decays:

$$\frac{1}{\Gamma} \frac{d\Gamma}{d\cos\theta_f} = \frac{1}{2} \left(1 + \mathcal{P}_t \kappa_f \cos\theta_f \right)$$

• 'spin analyzing power' κ_f depends on particle identity:

$$\circ \kappa_W = -\kappa_b$$

$$\kappa_{\ell} = 1, \qquad \kappa_{\nu} \simeq -0.3$$

- Hadronic tops: cannot distinguish \bar{d} from u
 - but: on average, \bar{d} less energetic than u in top rest frame
 - $\circ \Rightarrow$ form angular distribution of the softer light quark jet j
 - $\circ \kappa_i \simeq 0.5$ (Jezabek, 1994)

Hadronic tops

- Polarization-sensitive kinematic distributions for hadronic tops:
 - take axis of polarization to be top direction of motion in the lab frame: net polarization $\langle \mathcal{P}_D \rangle$

$$\frac{1}{\Gamma} \frac{d\Gamma}{d\cos\theta_i} = \frac{1}{2} \left(1 + \langle \mathcal{P}_D \rangle \kappa_i \cos\theta_i \right)$$

- o b-jet angular distribution: $\kappa_b \simeq -0.4$
- reconstructed W: $\kappa_W \simeq 0.4$
- o softer light quark jet: $\kappa_i \simeq 0.5$
- Depending on the methods of event selection and top reconstruction in a given study, any or all of these distributions may be appropriate.

Collimated tops

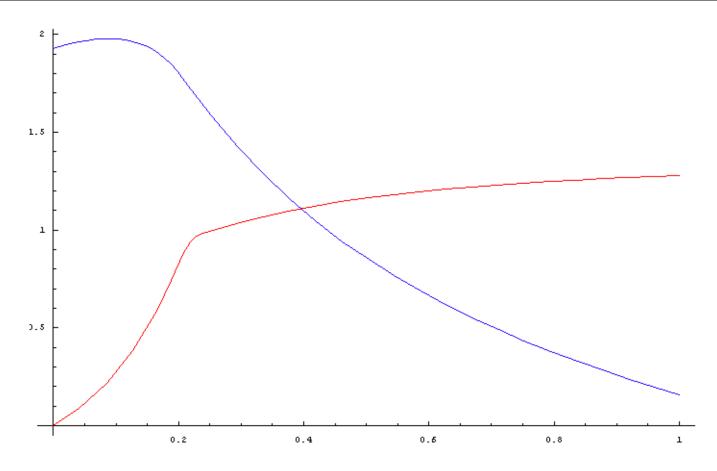
- Tops from very massive new physics will be highly boosted, and their decay products will be collimated
- Hadronic tops: if b (or W) can be identified within top, then a natural variable is the fraction of lab frame energy carried by the b, $z = \mathcal{E}_b/\mathcal{E}_t$
 - In collinear limit,

$$\frac{1}{\Gamma} \frac{d\Gamma}{dz} = \frac{m_t^2}{m_t^2 - m_W^2} (1 - \kappa_b \langle \mathcal{P}_D \rangle) + \kappa_b \langle \mathcal{P}_D \rangle \frac{2m_t^2}{m_t^2 - m_W^2} z$$

- o axis of polarization again top direction of motion in the lab frame
- Leptonic tops: natural variable is fraction of visible lab frame energy carried by the lepton,

$$u = \frac{\mathcal{E}_{\ell}}{\mathcal{E}_{\ell} + \mathcal{E}_{b}}$$

Collimated tops



Distribution of u for positive (red) and negative (blue) helicity tops.

• Measurable distribution a linear combination of positive and negative helicity distributions according to degree of polarization $\langle \mathcal{P}_D \rangle$

Polarization from cascade decays

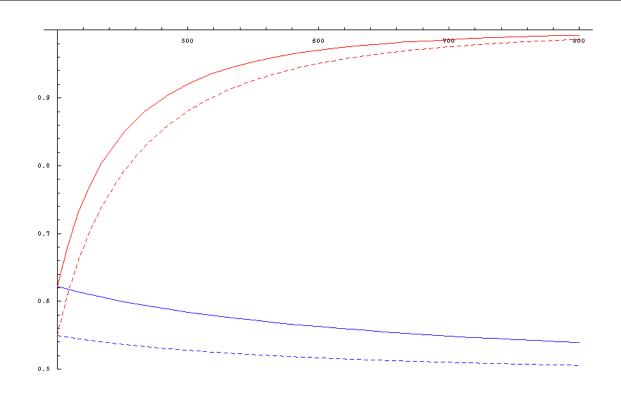
- Many well-motivated models feature cascade decays of the form (top partner) $\rightarrow t + (\text{vector boson partner})$
 - decay produces polarized tops due to chiral couplings λ_L, λ_R
 - can compute net production polarization $\langle \mathcal{P}_P \rangle$ as a function of particle masses
- SUSY: $\tilde{t} \to t \chi^0$

$$\langle \mathcal{P}_P \rangle = \frac{(|\lambda_R|^2 - |\lambda_L|^2)M|p_t|}{(|\lambda_L|^2 + |\lambda_R|^2)(ME_t - m_t^2) + 2\text{Re}(\lambda_L\bar{\lambda}_R)m_\chi m_t}$$

• Same-spin partner models: $\hat{T} \rightarrow t \, \hat{V}^{\mu}$

$$\langle \mathcal{P}_P \rangle = \frac{(|\lambda_R|^2 - |\lambda_L|^2)|p_t|(M^2 + 2m_V^2 - m_t^2)}{(|\lambda_L|^2 + |\lambda_R|^2)(3E_t m_V^2 + 2M|p_t|^2) - 6\text{Re}(\lambda_L \bar{\lambda}_R) m_V^2 m_t}$$

Polarization from cascade decays

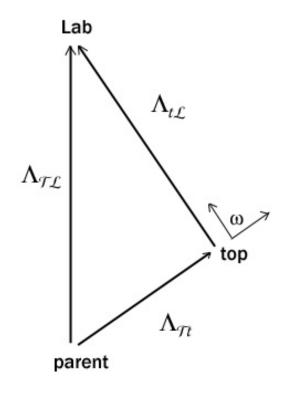


Net top quark polarization $\langle \mathcal{P}_P \rangle$ for purely chiral couplings (b=1, a=0).

- Red: fixed vector boson partner mass of 200 GeV
- Blue: vector boson partner mass of (top partner mass)—200 GeV
- o Solid (dashed) lines: opposite- (same-) spin partners

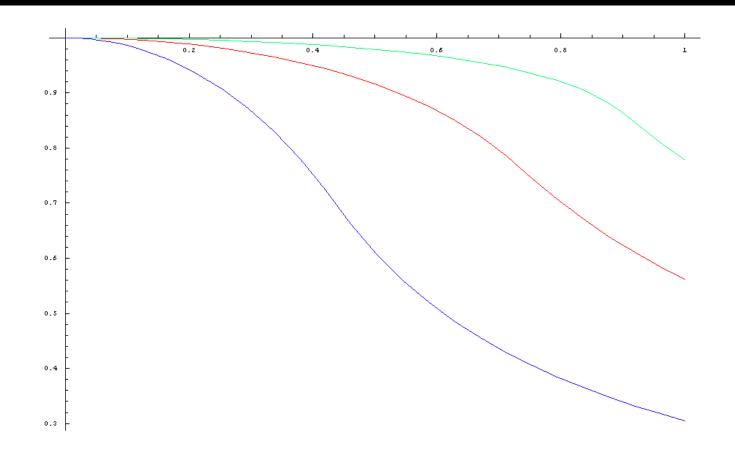
From production to detection

• Cascades give net polarization $\langle \mathcal{P}_P \rangle$ along direction of motion in the *parent* rest frame. Not directly observable!



- Observable polarization is reduced: $\langle \mathcal{P}_D \rangle = \langle \mathcal{P}_P \rangle \cos \omega$
- $\circ \cos \omega$: calculable function of particle masses, production dynamics

Wigner suppression



- $\cos \omega$ as a function of parent top partner boost β
 - Red: $(M_T, M_V) = (500 \text{ GeV}, 150 \text{ GeV})$
 - Green: $(M_T, M_V) = (900 \text{ GeV}, 300 \text{ GeV})$
 - Blue: $(M_T, M_V) = (900 \text{ GeV}, 700 \text{ GeV})$

Conclusions

- Top polarization is a valuable window into the chiral structure of new physics
 - finite m_t complicates the story, but detailed quantitative predictions are possible
- Polarization measurements in hadronic tops can be useful, especially if new physics contains stable invisible particles and decays through complicated cascades
- Also interesting are polarization-sensitive variables for collimated tops
 - here, two simple variables; others possible
 - \circ if b, W not separately discernible in collimated hadronic tops than more sophisticated techniques are necessary